3.2
The Amorphous State

Glasses

- **reversible response to small deformation** (no permanent stresses during a shear strain - unlike a liquid)
- **no long-range organisation**
- **high viscosity:** $\eta > 10^{11}\ Pa\ s$ (traditional definition)

coal (amorphous carbon)

silicate glasses

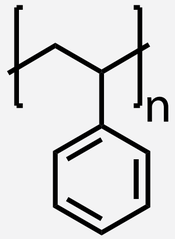
caramel

polymers

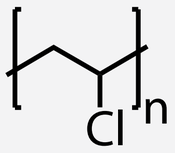
- **a glass is an amorphous and rigid solid; polymer glasses are often transparent.**

Amorphous, Glassy Polymers

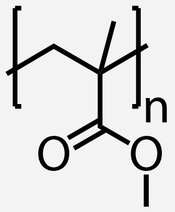
atactic polymers



PS

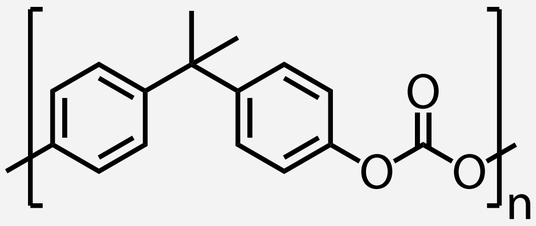


PVC

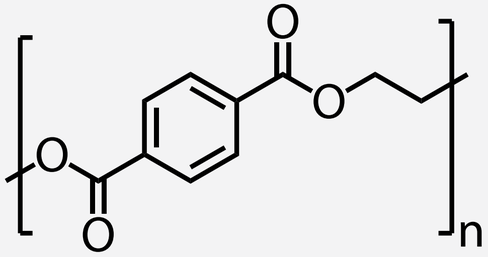


PMMA

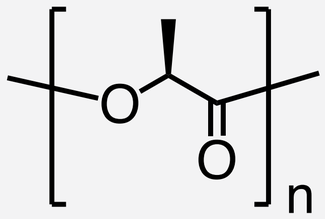
very slow crystallization



polycarbonate

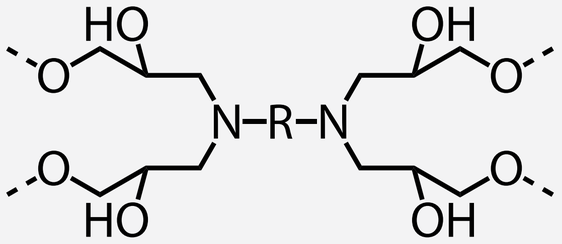


PET



PLLA

high cross-link density

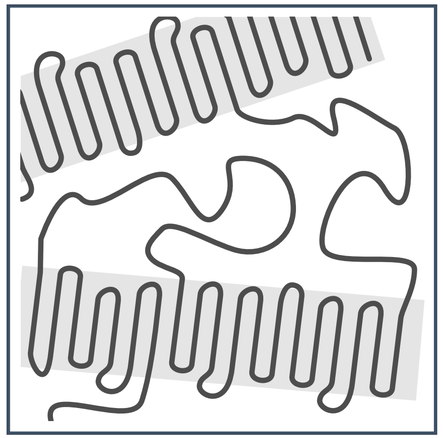


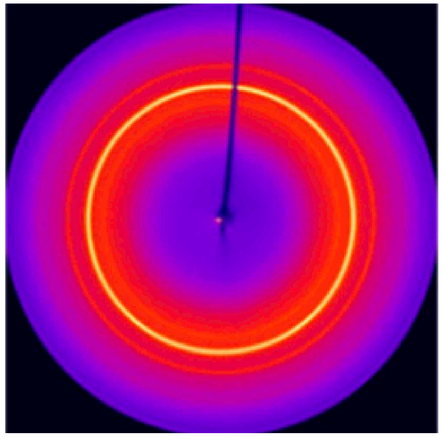
…plus various thermosets (e.g. epoxides)

The Amorphous State

- **random distribution of polymer chains in the matrix (no ordered structures)**
- **diffuse intra- and interchain peaks in X-ray diffraction patterns**

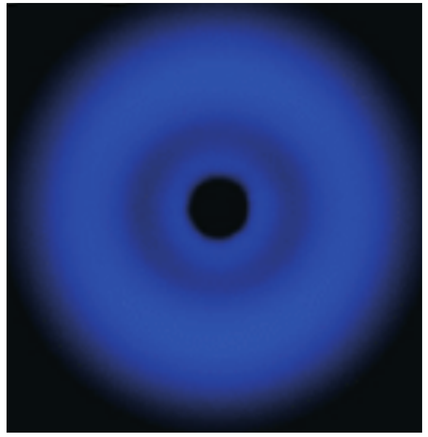
crystalline (sharp peaks)

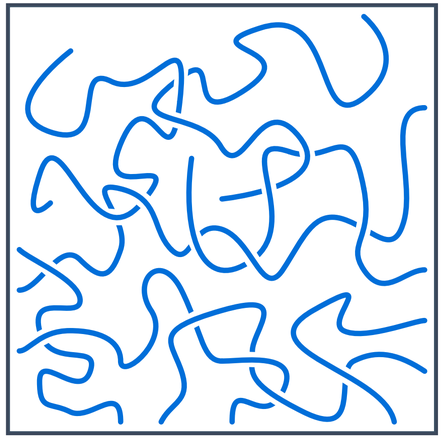




(isotropic sample)

amorphous (blurry peaks)

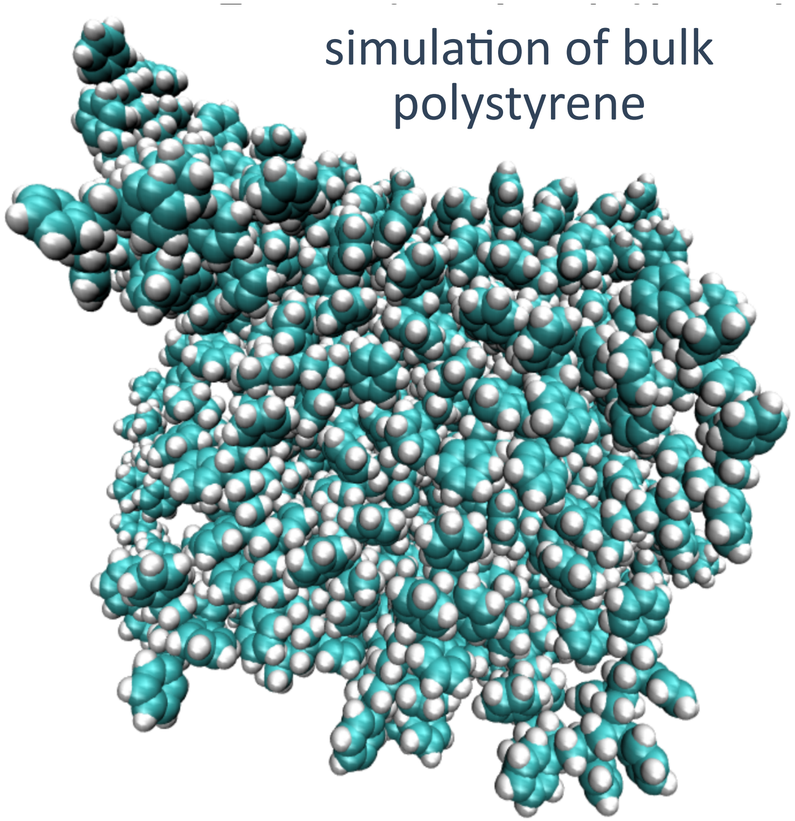


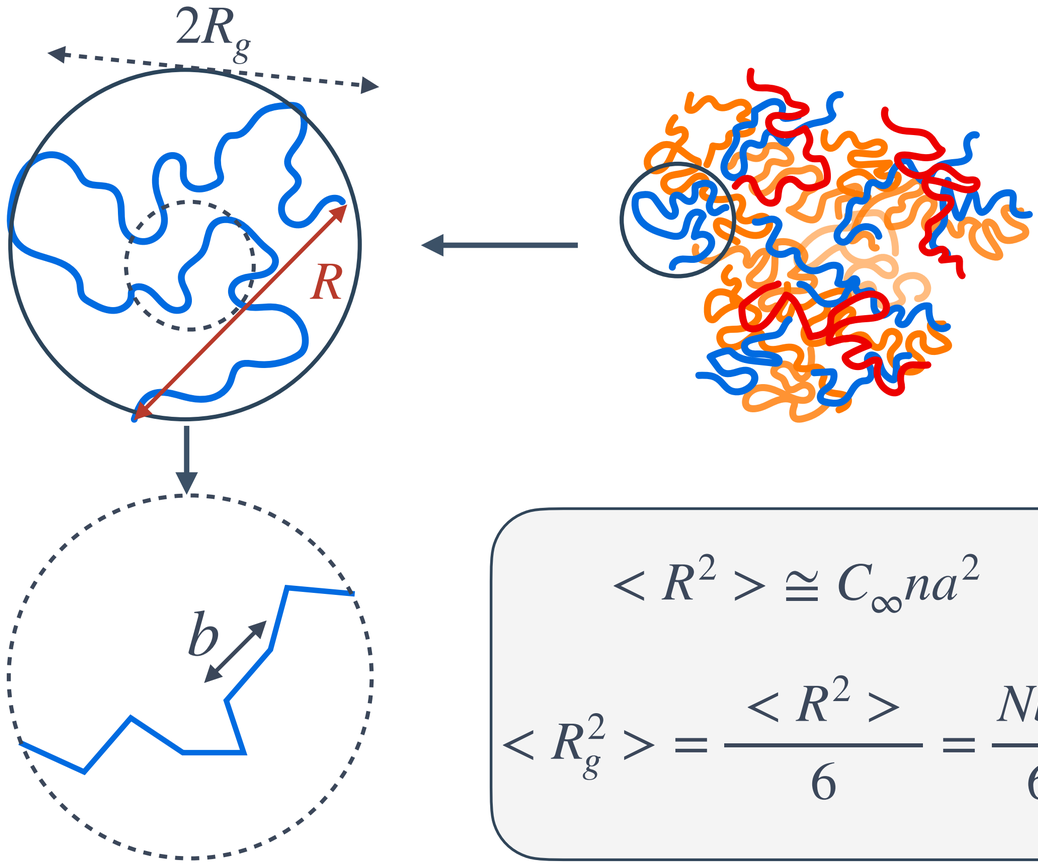


(isotropic sample)

- **local order does not exceed 1-2 nm in the molten state**
  (ample evidence from wide angle X-ray diffraction, light scattering, Brillouin scattering, Raman spectroscopy, …)

Conformations in the Amorphous State





$$< R^2 > \cong C_\infty n a^2$$

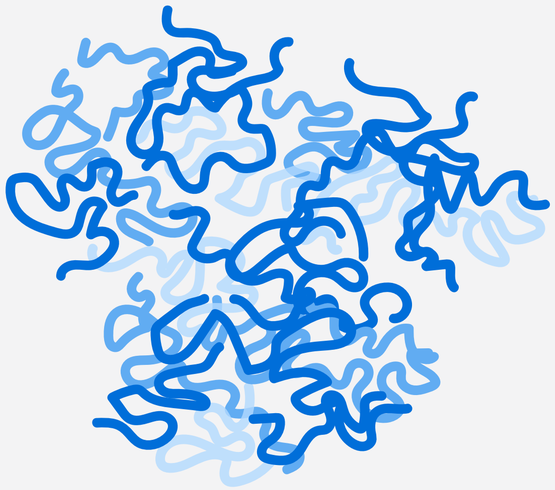
$$< R_g^2 > = \frac{< R^2 >}{6} = \frac{N b^2}{6}$$

- **chains adopt “random coil” conformation of an ideal isolated chain, if chain mobility is sufficiently high**
  (experimental evidence from neutron scattering and numerical simulations)

The Glass Transition Temperature

- **T_g: the temperature below which the conformations stop changing (on a large scale)**
- conformations in the glassy state are “frozen” ideal chain conformations of the molten state

glassy state



frozen conformations

rigid solid

T_g

rubbery or melt state

rapid conformational changes

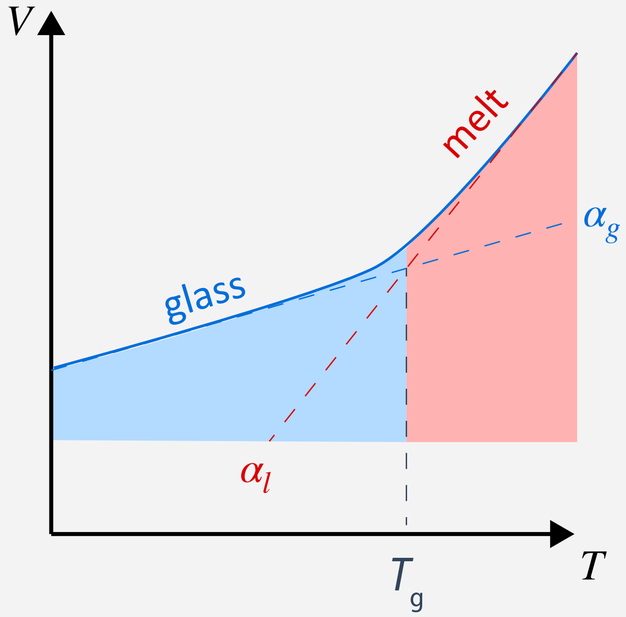
dynamic, mobile, fluidic

- **techniques for measuring T_g are based on properties that reflect changes in mobility**

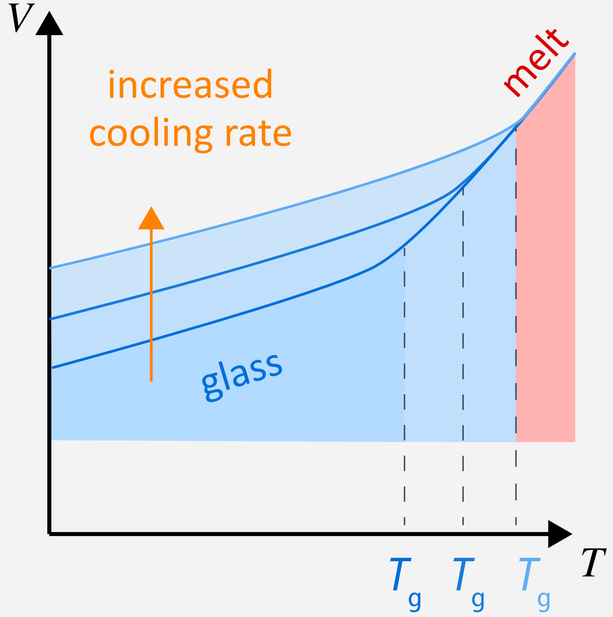
The Glass Transition Temperature

- **at T_g, sudden change in thermal expansion coefficient, α, upon cooling from the liquid state**
  (measured via a dilatometer)

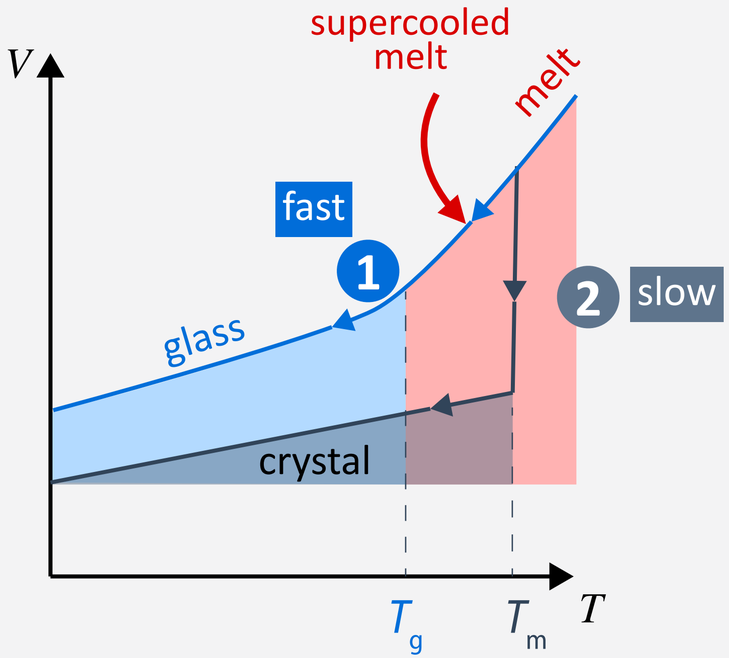
differences in α



cooling rate dependence



suppressed crystallization

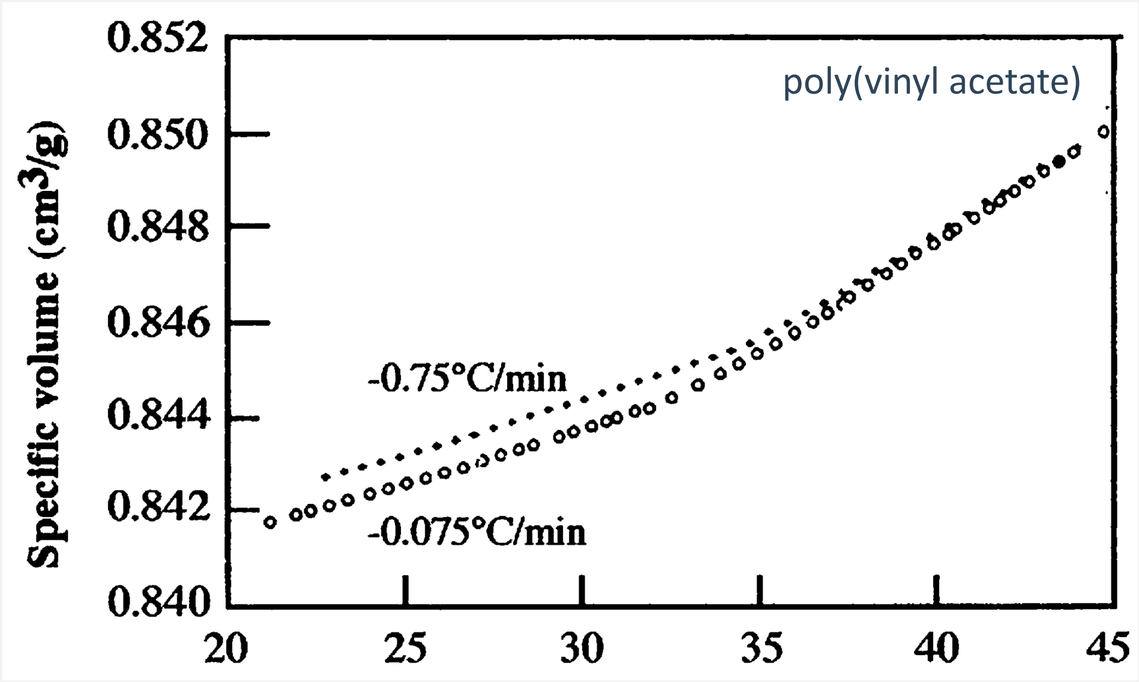


- **the transition temperature decreases with decreasing cooling rate**
- **T_g does not reflect thermodynamic equilibrium conditions! The glassy state is out-of-equilibrium!**

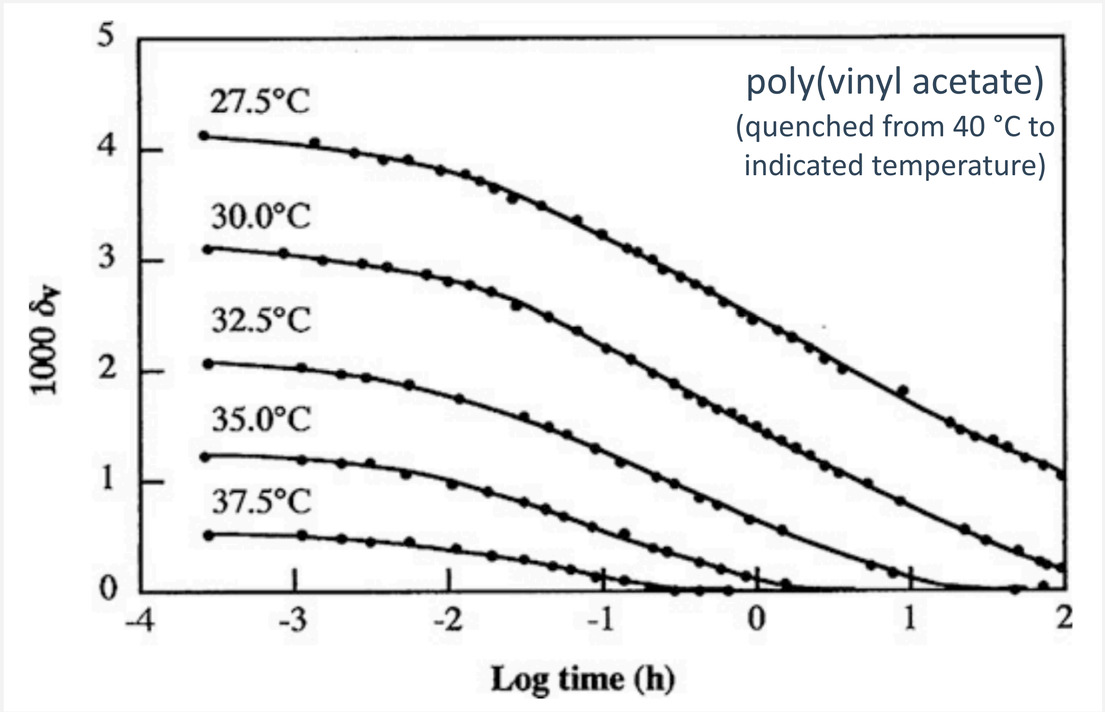
Physical Ageing

- isothermal volume recovery towards establishment of an equilibrium on long time periods

specific volume



volume recovery δ_v



- **physical ageing as a proof of the out-of-equilibrium nature of the glassy state**

Effect of Measurement Speed

- **T_g may vary considerably depending on measurement type and its typical standard condition**
- example: polystyrene

measurement Type	condition (rate)	T_g °C
electrical tests	1000 Hz	121
mechanical characterization	100 Hz	104
DSC	10 °C/min	100
dilatation (cooling)	2 °C/min	96
dilatation (cooling)	7 x 10^{-4} °C/min	82

- **the T_g is a fictive temperature and not a fundamental thermodynamic quantity to be determined**

1st vs. 2nd Order Thermodynamic Transitions

$$dG = VdP - SdT$$

1st partial derivatives of G:

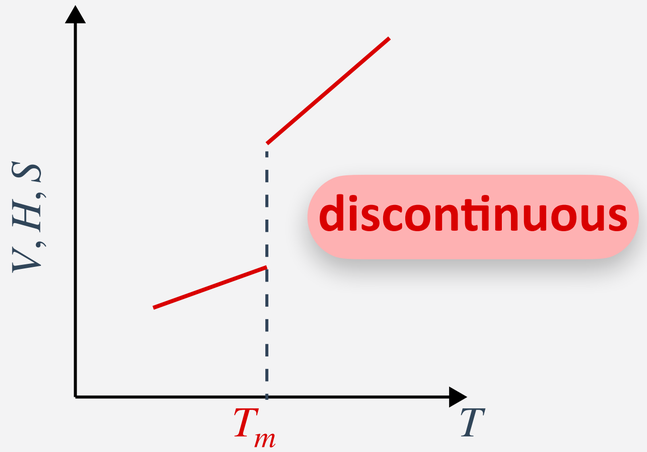
$$V = \left(\frac{\partial G}{\partial p}\right)_T \qquad -S = \left(\frac{\partial G}{\partial T}\right)_p$$

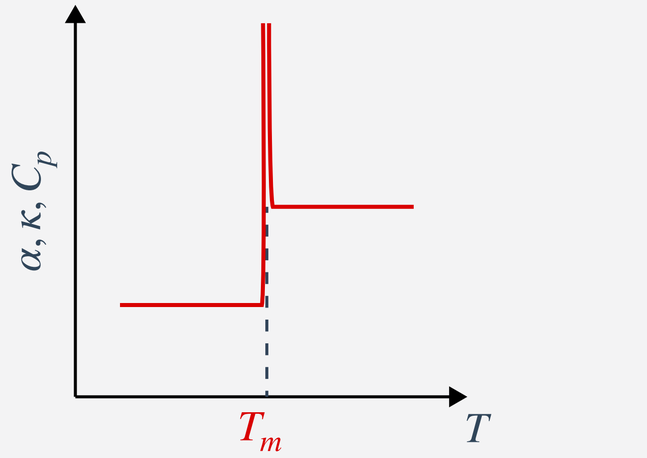
2nd partial derivatives of G:

$$\alpha = \frac{1}{V}\left(\frac{\partial V}{\partial T}\right) = \frac{1}{V}\left(\frac{\partial^2 G}{\partial T \partial p}\right)$$

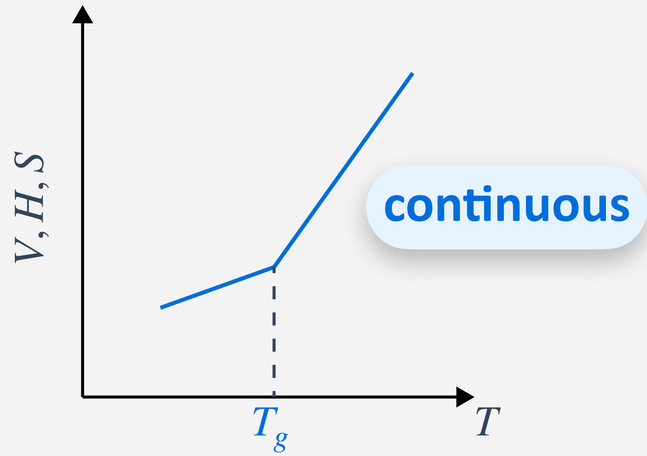
$$C_p = T\left(\frac{\partial S}{\partial T}\right) = -T\left(\frac{\partial^2 G}{\partial T^2}\right)$$

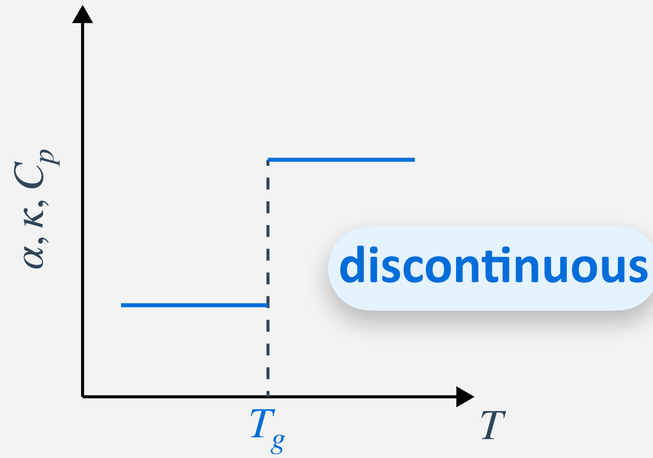
1st order transition
(polymer crystallisation)





2nd order transition
(glass transition)

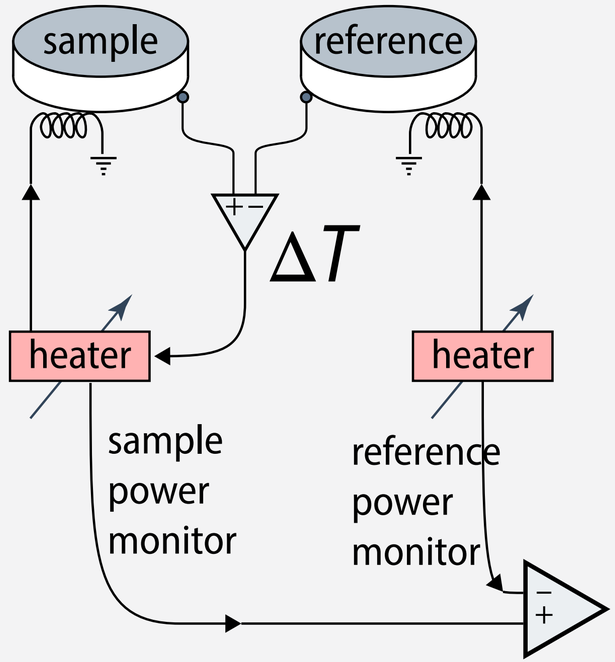


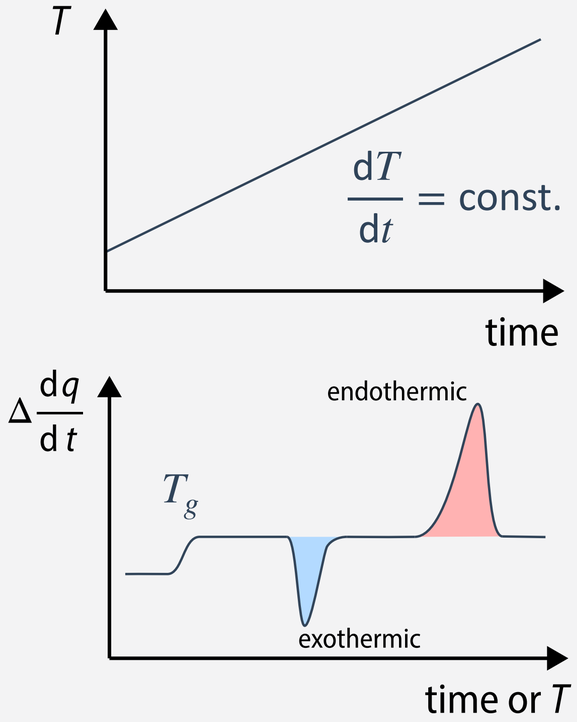


- **glass transitions: characteristics of a 2nd order transition, but NOT a true thermodynamic transition**

Principles of Differential Scanning Calorimetry

- Differential Scanning Calorimetry (DSC) monitors heat effects associated with phase transitions





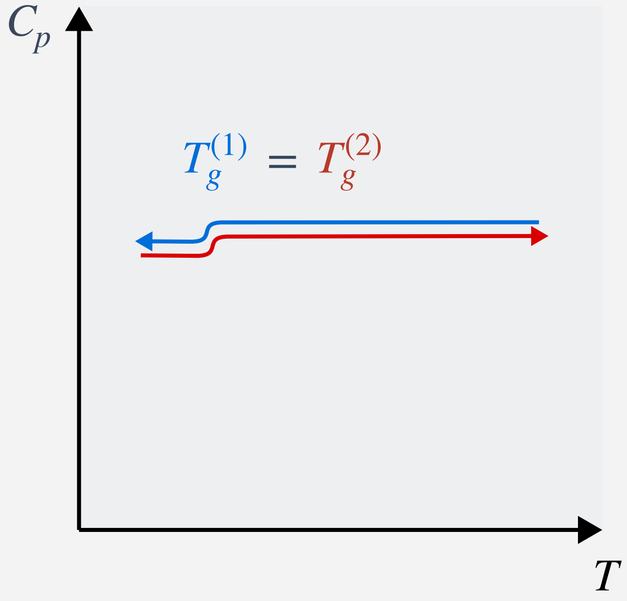
$$C_p = \frac{dQ}{dT} = \frac{dH}{dT}$$

- **endothermic process (e.g. polymer melting): heat is absorbed by the sample ($\Delta \mathrm{d}H/\mathrm{d}t > 0$)**
- **exothermic process (e.g. polymer crystallization): heat is released by the sample ($\Delta \mathrm{d}H/\mathrm{d}t < 0$)**
- **glass transitions cause an endothermic baseline shift**

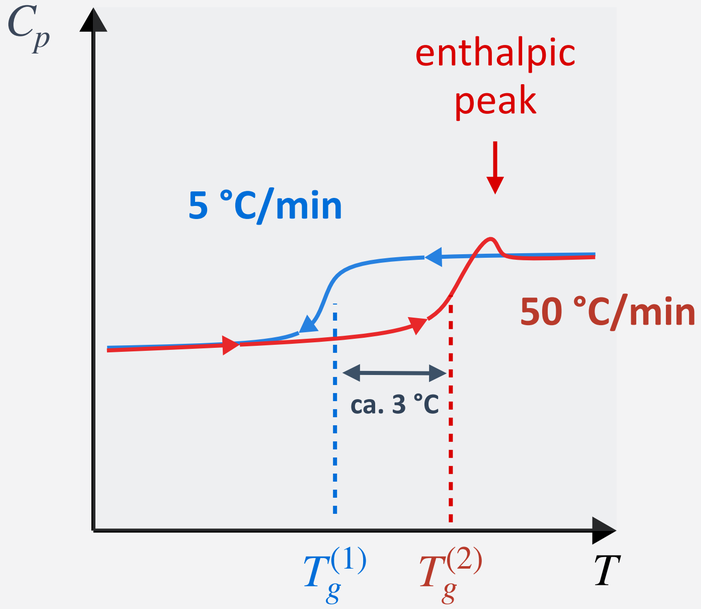
Heat Capacity Changes in Differential Scanning Calorimetry

- **at T_g, sharp drop in thermal capacity upon cooling from the liquid state:** $C_p = \frac{dQ}{dT} = \frac{dH}{dT}$
- **during heating, the transition is often accompanied by an enthalpy peak**
  (an indication of the metastable nature of the glassy state)

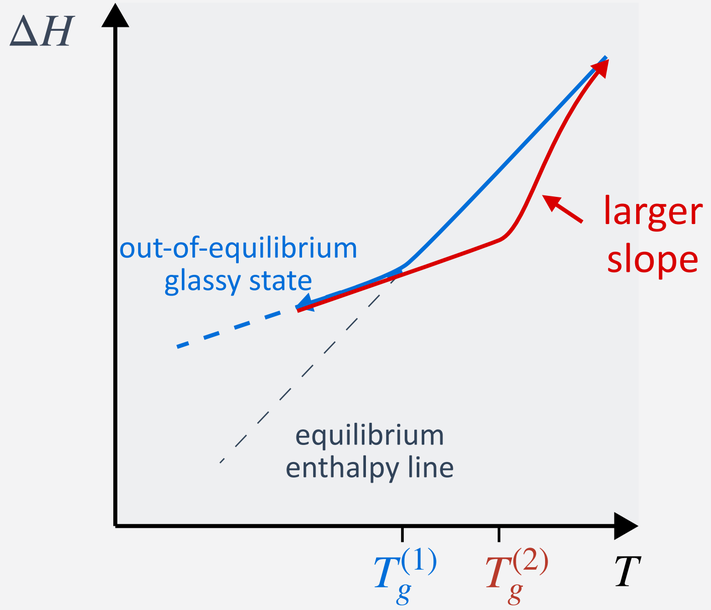
same cooling and heating rate



different cooling and heating rate



enthalpy changes

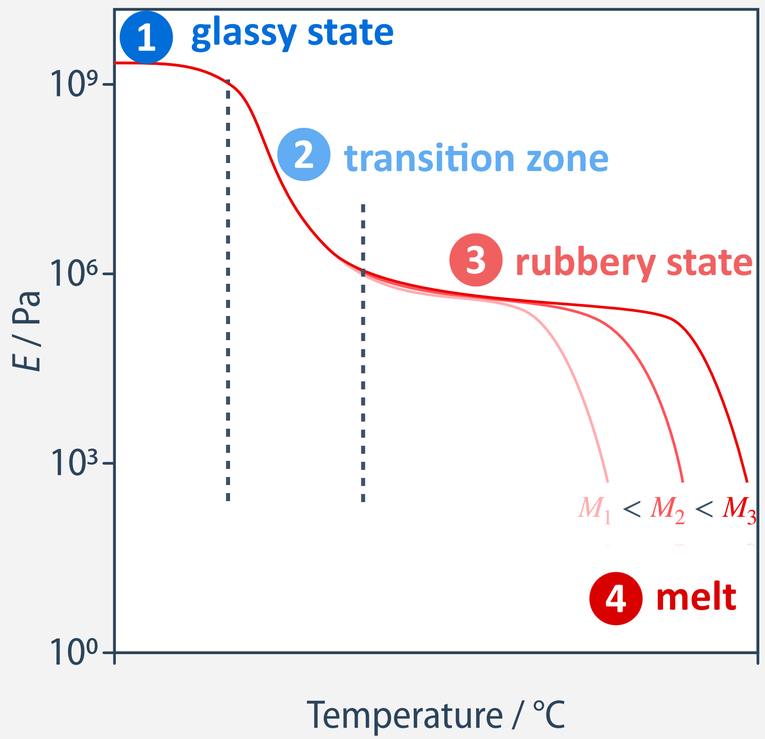


- **glassy state is out of equilibrium: chains cannot change conformation to establish equilibrium at a given *T***

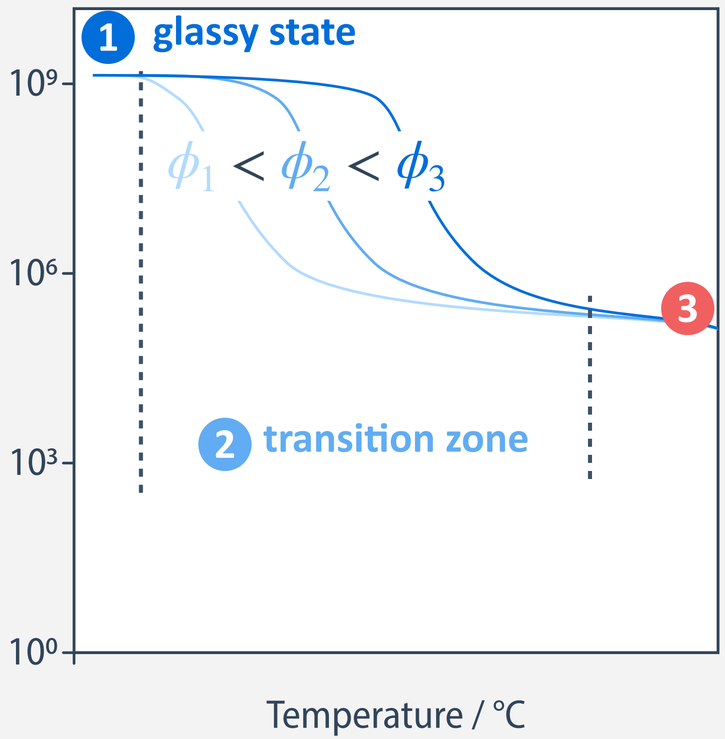
Mechanical Properties

- **at T_g, sharp increase of shear modulus** (+ other elastic moduli!) **from 10^6 to 10^9 Pa**
  (upon cooling from melt state)

universal trend of Young's modulus



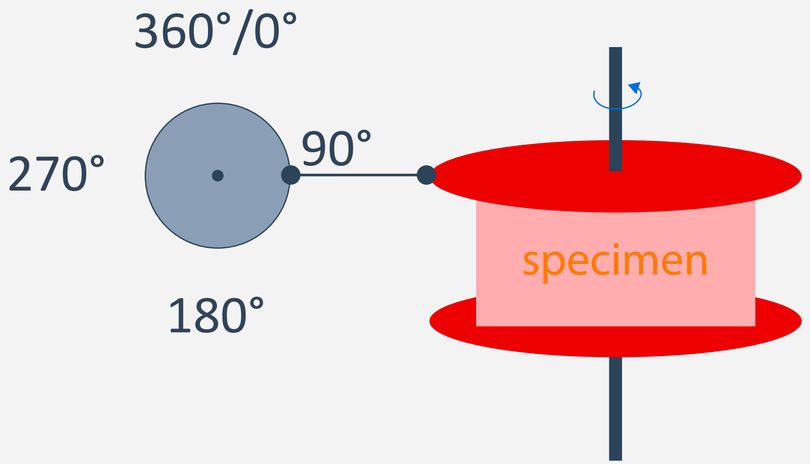
cooling rate dependence



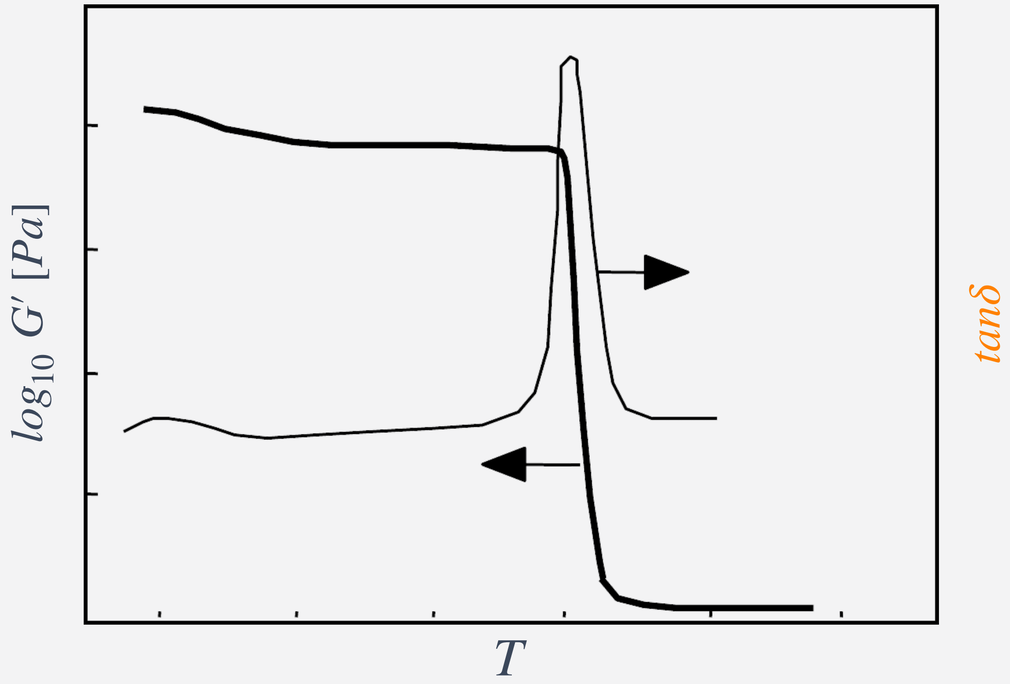
- **for long polymer chains, a rubbery regime is displayed at temperatures immediately above T_g**
  (the extent of this rubbery state depends on the molecular weight! see **Chapters 4.1 and 4.2**)

Dynamic Mechanical Analysis / Shear Rheology

- application of small deformations in a cyclic manner allow to study the materials response to stress, temperature, and frequency
- information on damping ($tan\ \delta$): a measure of energy dissipation
- at T_g, a characteristic damping peak is observed



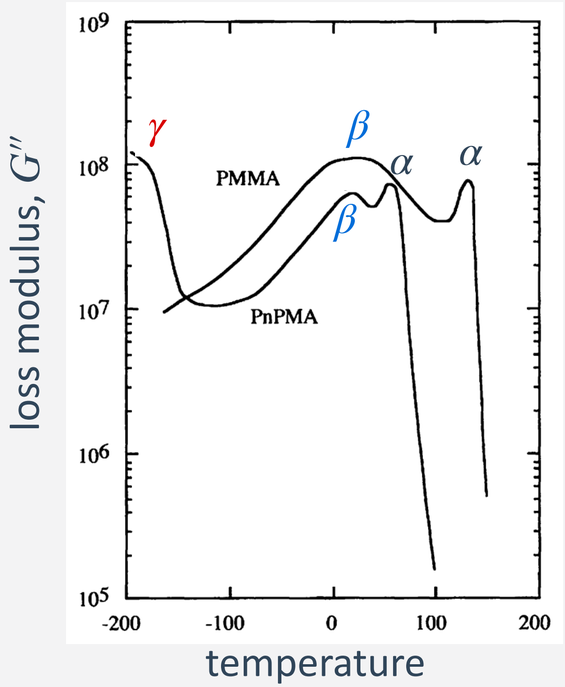
for working principle: see **Chapter 4.2**

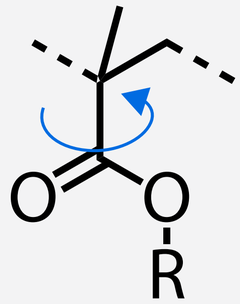


- T_g increases with increasing frequency (smaller experimental time scale).

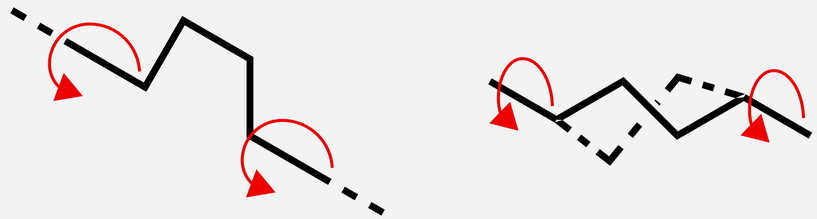
Relaxation Processes in the Glassy State

- sudden change in physical properties at the glass transition temperature such as the loss modulus (G'')
- small subunits remain mobile below T_g and therefore mechanically active towards lower temperatures





rotation



crankshaft motion

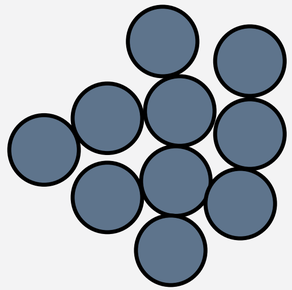
- these energy dissipation mechanism are unique to polymers, and not encountered in inorganic glasses
- origin of plasticity: polymer glasses are still tougher than ceramics! (see Chapter 4.3)

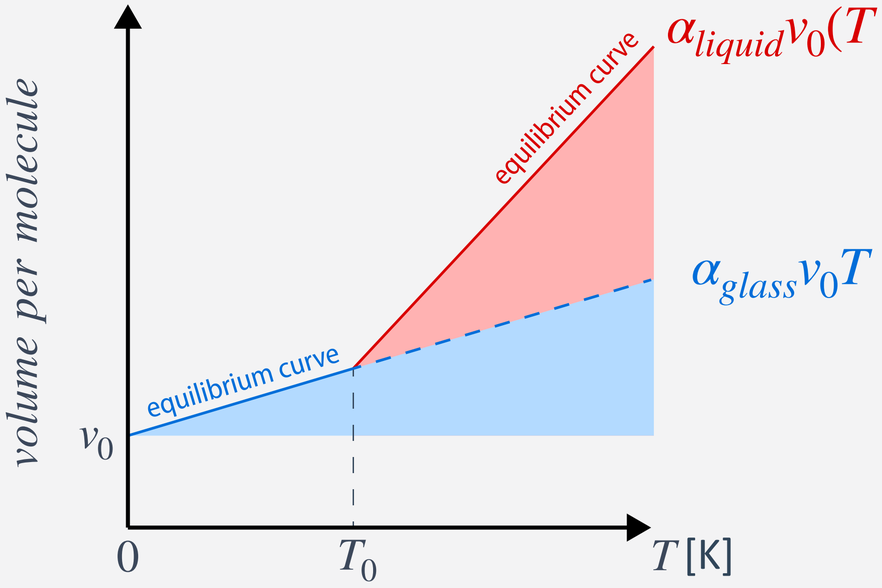
The Theory of Free Volume

- **at T_0, part of the volume created by thermal vibrations is "free", leading to holes of size v_f.**
- **movement, if locally $v_f \geq v_0$** (v_0: the volume occupied by a sphere/molecule)

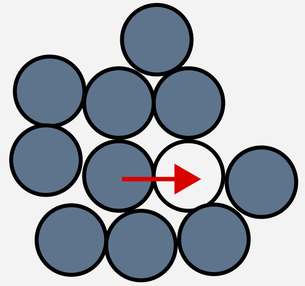
infinitely slow cooling

conformational changes restricted





conformational changes possible



$v_f \geq v_0$

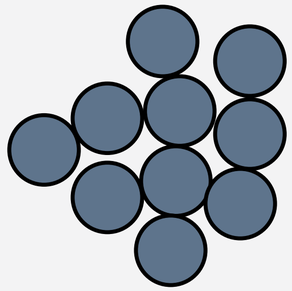
- **no translational, rotational, and conformational movement possible below T_0**

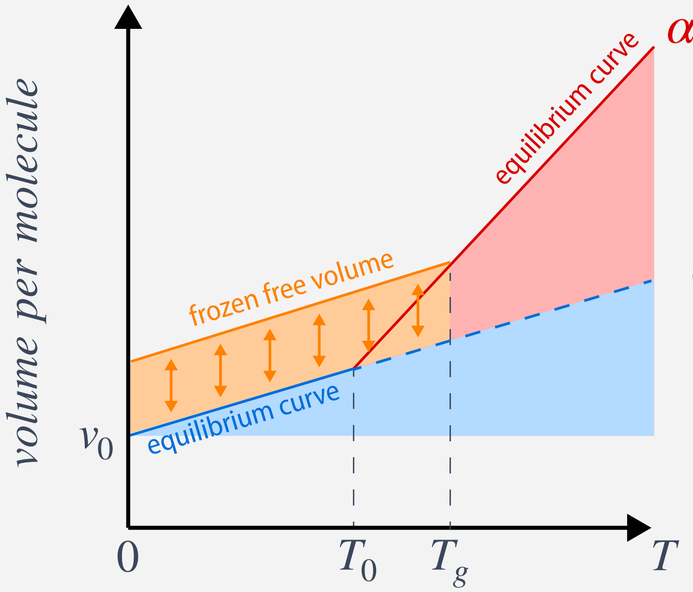
The Theory of Free Volume

- an infinitely slow cooling rate can in practice not be realized and T_0 cannot be measured
- free volume trapped in the glassy state (see exercise): $v_{fm} = v_0(T - T_0)\Delta\alpha$

finite cooling rate

conformational changes restricted

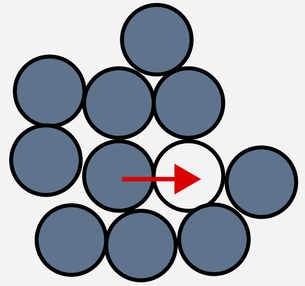




$\alpha_{liquid} v_0(T - T_0) + \alpha_{glass} v_0 T_0$

$\alpha_{glass} v_0 T$

conformational changes possible



$v_f \geq v_0$

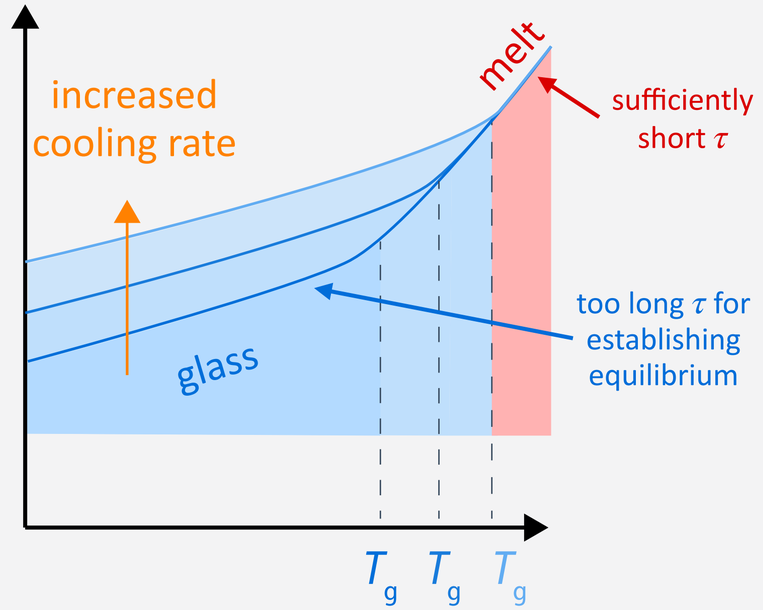
- at a finite cooling rate, holes become "frozen" at a temperature $T_g > T_0$, which depends on the cooling rate

Relaxation Times and Free Volume

- **relaxation time τ needed to change conformation for equilibration**
- free volume interpretation for quantification of changes in dynamic and mechanical properties:

relaxation time τ : $\tau = \tau_0 e^{\frac{v_0}{v_{fm}}} = \tau_0 e^{\Delta\alpha^{-1}/(T-T_0)}$ $\Delta\alpha = \alpha_{liuqid} - \alpha_{glass}$

viscosity η (proportional to τ): $\eta = \eta_0 e^{\frac{v_0}{v_{fm}}} = \eta_0 e^{\frac{\Delta\alpha^{-1}}{T-T_0}}$



- **empirically verified relationships for many glass forming systems**
- **limitations: free volume theory was founded for simple liquids** (macromolecular chain connectivity, thermal activation of chain movement, secondary interactions not considered)

Thermodynamic Theory by Gibbs-Di Marzio

- chain connectivity and thermal activation are taken into account
- decreased difference in entropy between liquid and solid phase upon supercooling

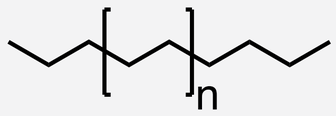
$$S_l - S_{glass} = \Delta S_g - \int_T^{T_g} \Delta C_p ln T$$

- Kauzmann paradoxon: $S_l < S_{glass} \; for \; T < T_g$
- a phase transition must take place between T_g and T_0
- there is an underlying thermodynamic basis for the glass transition, but experimentally never verified

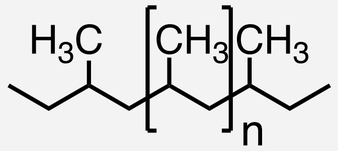
Structural Factors Impacting T_g

Chain Rigidity

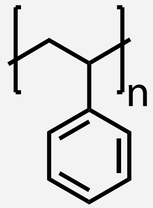
- decreased chain mobility at increased sterically demanding (bulky) side groups



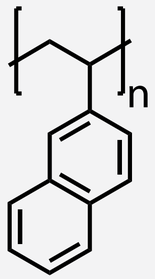
T_g = -100 °C



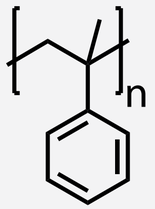
T_g = -10 °C



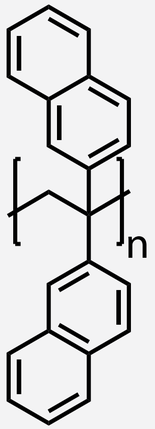
T_g = 100 °C



T_g = 135 °C



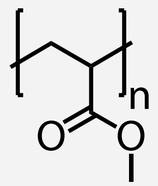
T_g = 175 °C



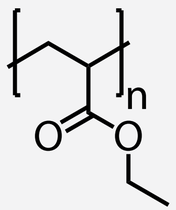
T_g = 264 °C

Plasticisation

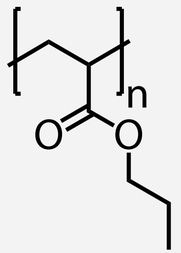
- internal plasticisation: addition of flexible substituents reduce the T_g
- external plasticisation: addition of small amounts of a solvent (mainly used for PVC)



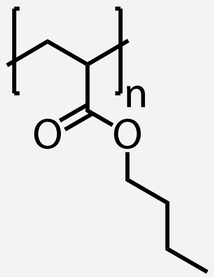
$T_g = 6$ °C



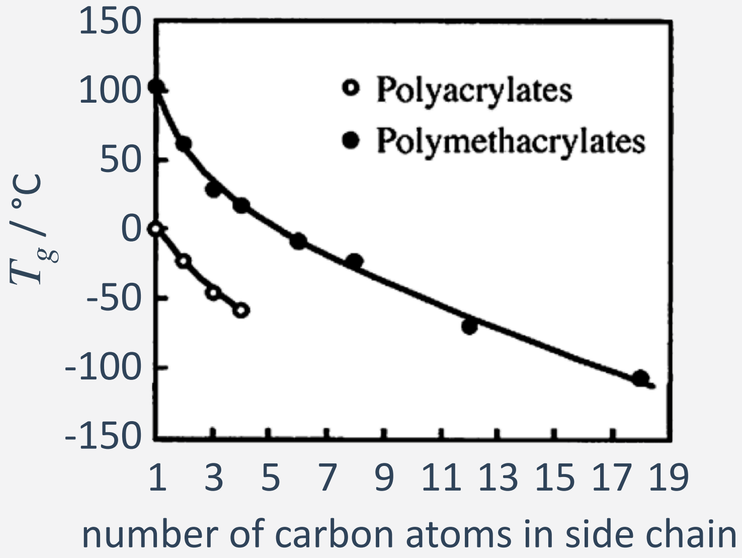
$T_g = -24$ °C



$T_g = -48$ °C



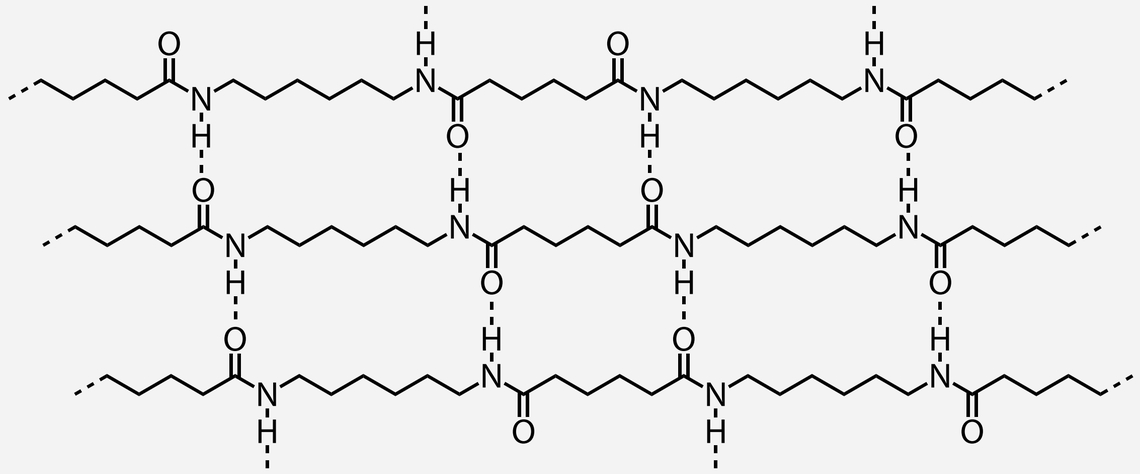
$T_g = -55$ °C

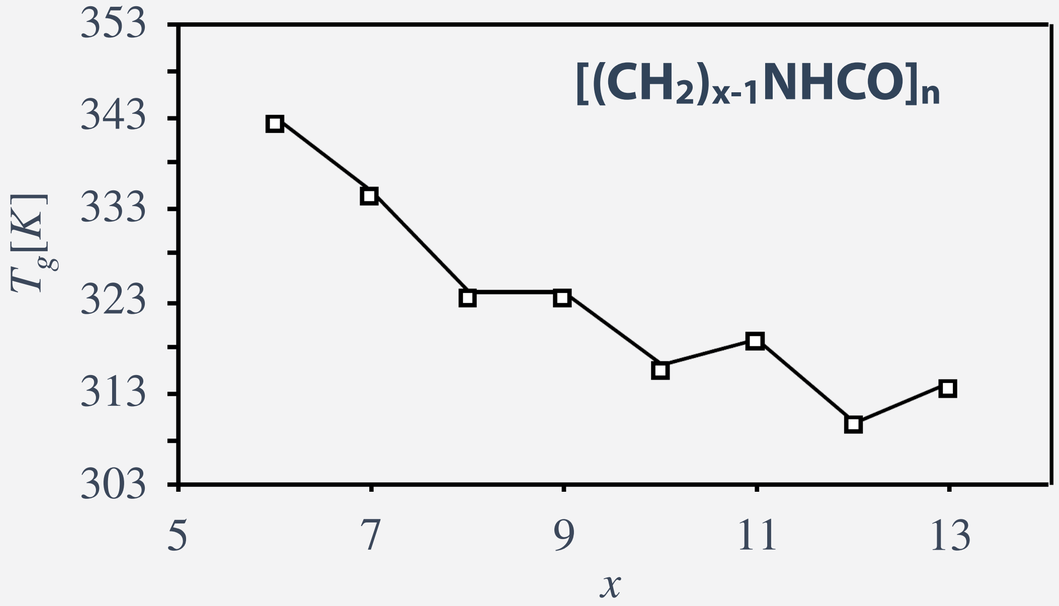


- self-plasticisation of certain polyolefin's due to a wide distribution of the molar mass
- antiplasticizing effect in case of strong and specific polymer-solvent interactions

Specific Interactions

- decreased chain mobility in case of strong inter-chain interactions, i.e. in presence of polar groups (CN, NH, C=O) and, in particular, hydrogen bonds.





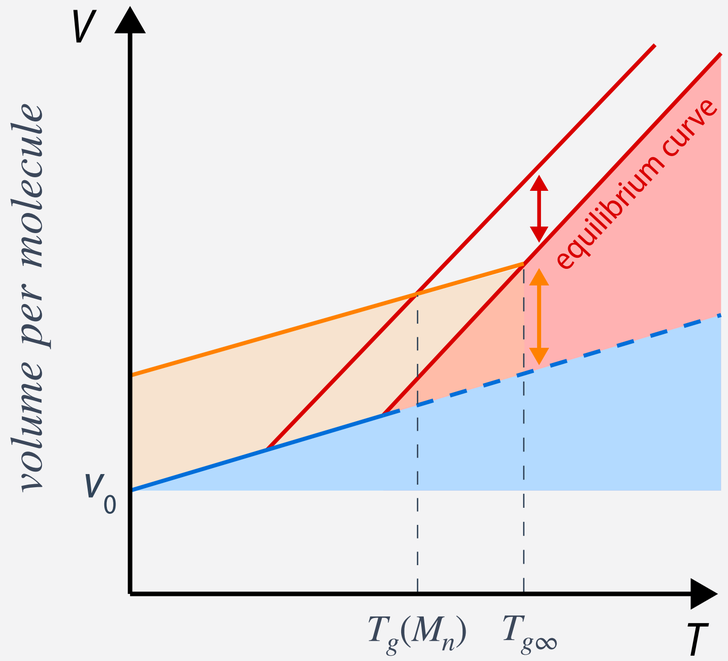
- increases in T_g at increased number of amide groups in nylons.
- solvents competing for hydrogen bonding interactions (i.e. H_2O) serve as plasticisers

Molecular Weight Dependence

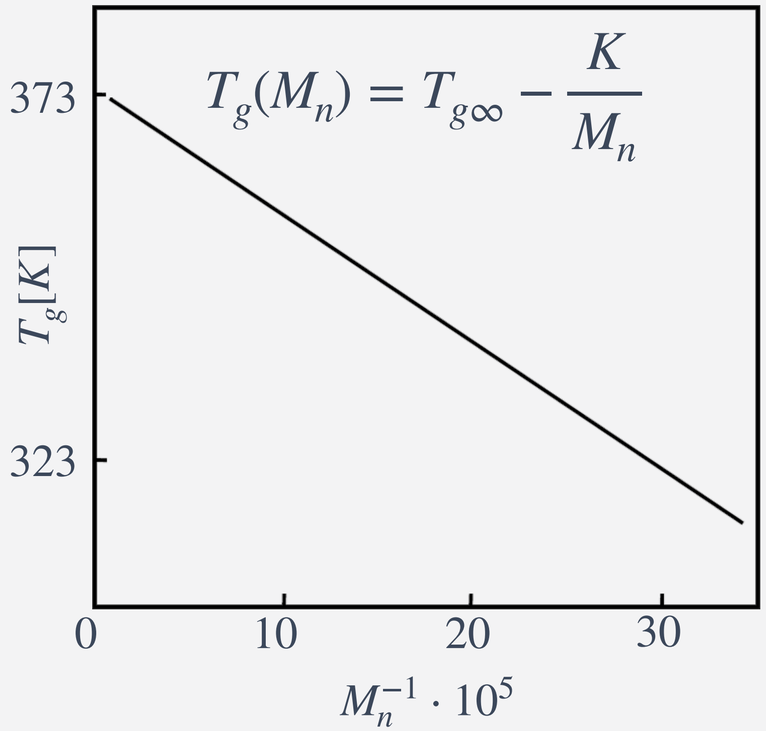
- the total fractional free volume due to chain ends, f_c, adds to the free volume:

$$v_{fm} = (\alpha_l - \alpha_g)v_0(T_{g\infty} - T_0) + f_c v_0 = (\alpha_l - \alpha_g)v_0(T_{g\infty} - T_0) + v_0\frac{2\rho N_A \theta}{M_n}$$

θ: additional free volume per end group



example: polystyrene



- **each chain end provides additional "free volume", θ (see also Exercise Sheet)**

Superglue

- working principle: strong adhesive nature of cyanoacrylates and quick polymerization (glass formation!)

The History of Superglue

1942

Cyanoacrylates discovered during WWII search gun sight plastics. Stick to everything, so discarded

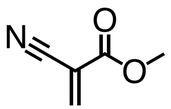
1951

Cyanoacrylates rediscovered during research looking for polymers for jet canopies

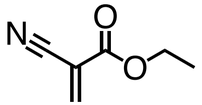
1958

Their potential finally realised: cyanoacrylates are developed into a glue which eventually becomes available commercially in 1958. Numerous other manufactures follow suit.

Cyanoacrylates



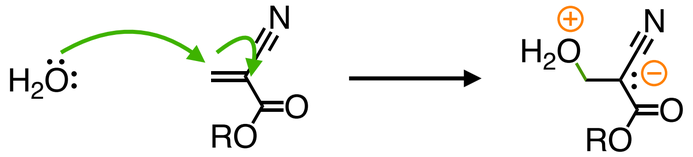
methyl cyanoacrylate



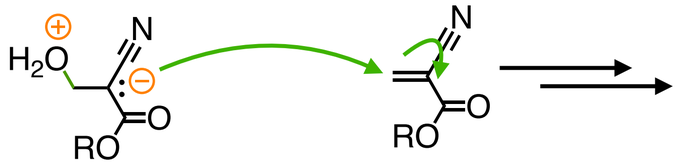
ethyl cyanoacrylate

Most common: ethyl cyanoacrylate, but others can also be used. Medical grade cyanoacrylates such as 2-octyl cyanoacrylate can be used to close up wounds

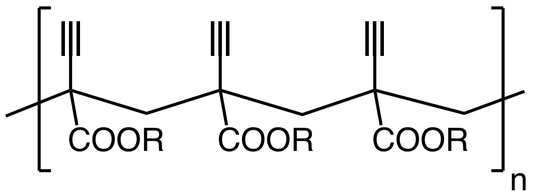
Mechanism



Cyanoacrylates 'cure' in the presence of water. Only a small amount of water is required to kick of the reaction - even the water vapour in air is enough.

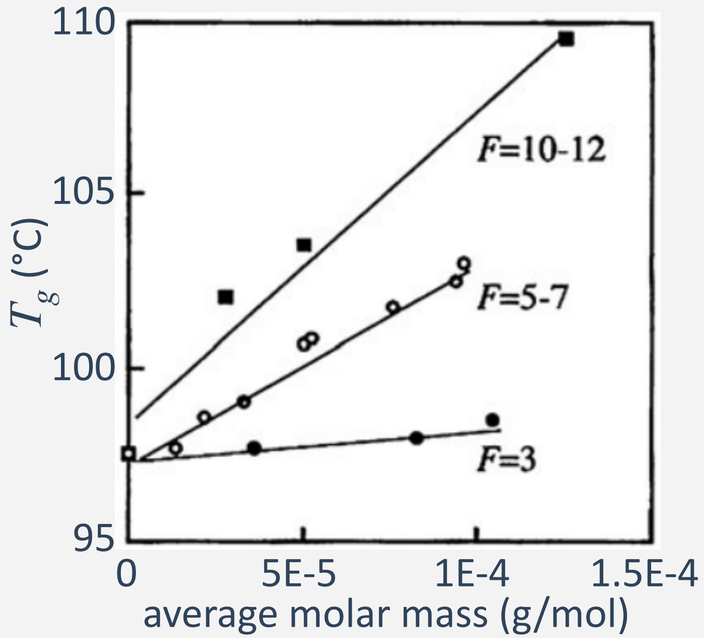


The reaction produces an anion which can add to more of the original cyanoacrylate, a process that repeats to form the adhesive polymer chains

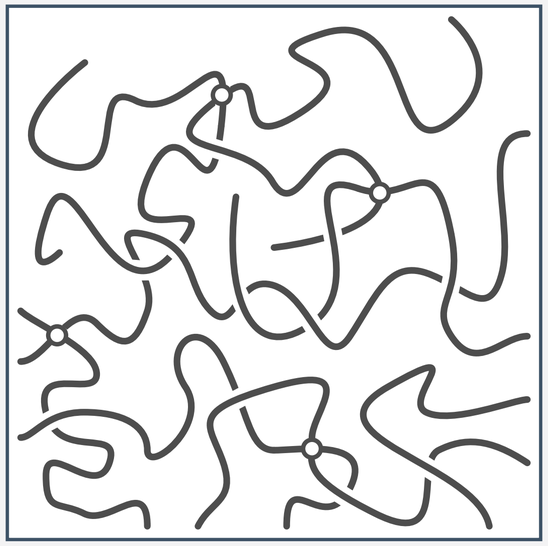


Cross-Linking

- **low cross-linking density: discussed relations hold true, in particular molecular weight dependence**
- **high cross-linking density: the sterical demand of catenary connections may lead to an increase in T_g**

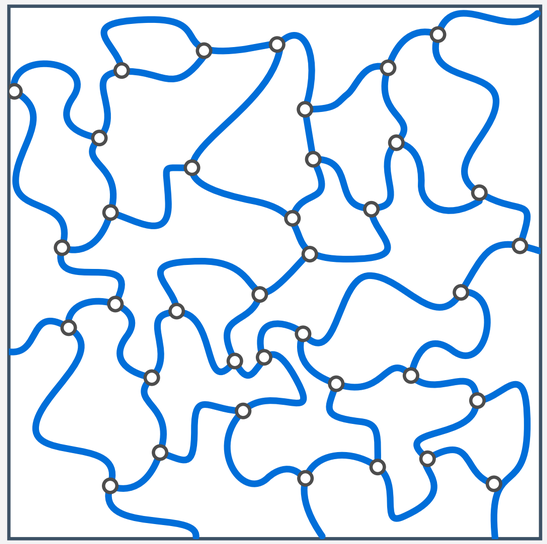


elastomers



$T_g <$ room temperature

thermosets



$T_g >$ room temperature

- **increased cross-linking density increases the T_g of thermosets**
- **impact: as thermoset fabrication proceeds, gradually decreased mobility can slow down synthesis**
  (hence, the need to cross-link at high *T* in some cases)

Copolymers

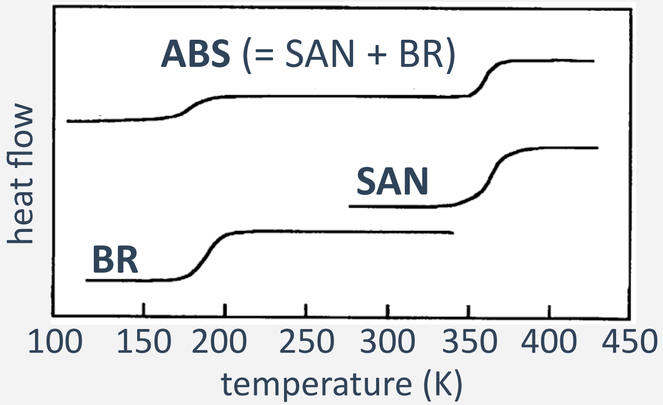
- **random copolymers, likewise homogeneous polymer mixtures, display a single T_g**

simple, empirical law of mixing:

$$\frac{1}{T_g} = \frac{w_1}{T_{g1}} + \frac{w_2}{T_{g2}}$$

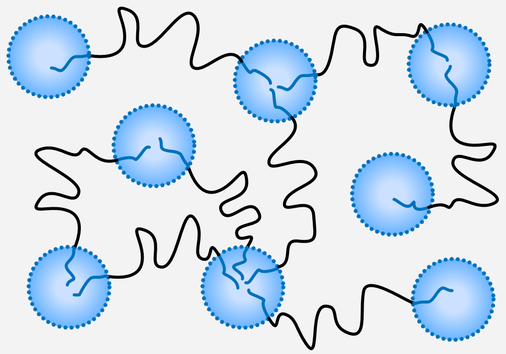
w_i: mass proportion of each component

- **several glass transitions in block copolymers depending on the number of components**



(also phase separated mixtures of polymers, **Chapter 5**)

- **thermoplastic elastomers: cross-links from glassy domains from one block at operating temperature**



glassy, hard

physical cross-links

∅ 10 – 100 nm

soft matrix: low T_g

Overview of Glass Transition Temperatures in Linear Polymers

Polymer Type $[CH_2\text{-}CHR]_n$	Substituent R	T_g [°C]	Trend Rational
polyethylene	H	-80	increasingly bulky substituents ↓
polypropylene	CH_3	-20	
polystyrene	Ph	100	
polybutene	C_2H_5	-24	**internal plasticisation:** increasingly flexible substituents ↓
polypentene	C_3H_7	-40	
Polyhexene	C_4H_9	-50	
poly(methylacrylate)	$COOCH_3$	5	**internal plasticisation**
poly(ethylacrylate)	$COOC_2H_5$	-20	
poly(propylacrylate)	$COOC_3H_7$	-48	
poly(butylacrylate)	$COOC_4H_9$	-55	
polyvinylalcohol	OH	85	**polar groups**
polyvinylchloride	Cl	81	
polyacrylonitrile	CN	105	
polyvinylacetate	$C(O)CH_3$	28	

Poly(methyl methacrylate) $[CH_2\text{-}CCH_3R]_n$	Substituent R	T_g [°C]	Trend Rational
PMMA (isotactic)	$COOCH_3$	145	**tacticity** for disubstituted repeat units
PMMA (atactic)	$COOCH_3$	105	
PMMA (syndiotactic)	$COOCH_3$	115	

Learning Outcome

- in the amorphous state, polymer chains adopt their ideal random walk conformations. These become “frozen in” at temperatures below the glass transition temperature
- determination of T_g by dilatometry, calorimetry, static mechanical tests, and dynamic measurements
- the out-of-equilibrium nature of the glassy state, like the strong dependence on measurement speed and provided experimental time can be rationalised by the theory of Free volume
- high T_g values are favoured by high M_n, chain stiffness, strong interchain forces (specific interactions)
- if necessary, plasticisers can be used to reduce the T_g